

TMD factorization and artemide

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ARTEMIDE

The only publicly-available code working within **TMD factorization**.

News: a lot of done both in theory and in phenomenology

- ▶ Critical revision and cross-checks of TMD factorization theory
- ▶ Complete theory justification of ζ -prescription used in **artemide**
- ▶ New perturbative calculations: NNLO (tw-2 gluon, polarized, etc.)
LO+NLO (tw-3)(!)
- ▶ Phenomenology: DY, π DY, SIDIS, DIS+jet.

artemide checklist

- ✓ ...
- ✓ Higgs production and gluon TMDs
- ✓ target-mass corrections for SIDIS
- ✓ joined fit of SIDIS+DY (\leftarrow we are here)
- power corrections for DY
- SSA and Sivers functions
- ...

TMD factorization

... should not be mistaken with resummation approach

- ▶ Resummation: $\Lambda_{QCD} \ll q_T \ll Q$
 - ▶ Common structure at $\Lambda_{QCD} \ll q_T \ll Q$ for some processes
(e.g. unpolarized DY and SIDIS, SSA).
- TMD factorization: $q_T \ll Q$

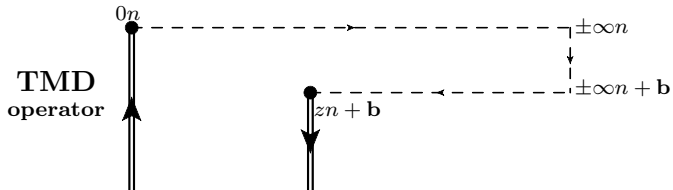
Pros:

- ▶ Can be applied down to low energies $Q \sim 2\text{GeV}$ and down to $q_T \sim 0$.
- ▶ Well established theory foundation (operators, renormalization, all that).
- ▶ Universal evolution
- ▶ Gives access to interesting non-perturbative functions.

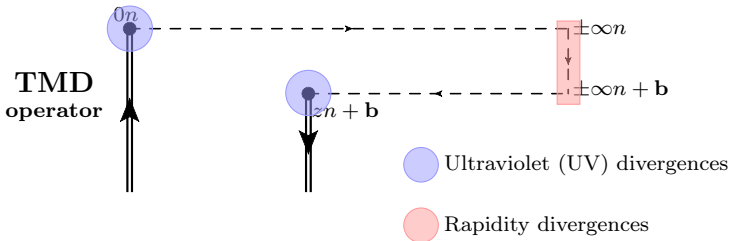
Cons:

- ▶ Plenty of new unknown functions.
- ▶ Only double-inclusive processes: $DY, SIDIS, ee \rightarrow hh + X, DIS + jet$.
- ▶ $q_T \ll Q$ and, in addition, q_T is not p_T (next slides)

TMD distributions have very peculiar operator structure



TMD distributions have very peculiar operator structure



TMD distribution depends on 2 scales

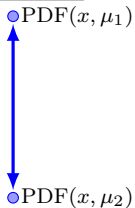
its evolution is given by 2 equations

$$\mu^2 \frac{dF(x, b; \mu, \zeta)}{d\mu^2} = \gamma_F(\mu, \zeta) F(x, b; \mu, \zeta),$$

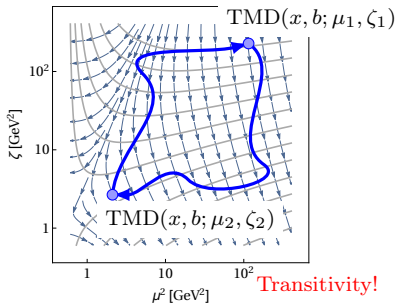
$$\zeta \frac{dF(x, b; \mu, \zeta)}{d\zeta} = -\mathcal{D}(\mu, b) F(x, b; \mu, \zeta)$$

Double-scale evolution rises a different set questions

1D evolution



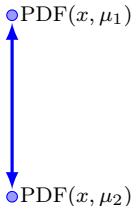
2D evolution



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Double-scale evolution rises a different set questions

1D evolution

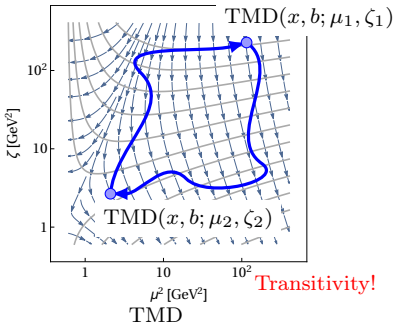


DIS

$$C(x, Q/\mu) \otimes q(x, \mu)$$

out: $\mu \sim Q$
in: $\mu \sim 2\text{GeV}$

2D evolution



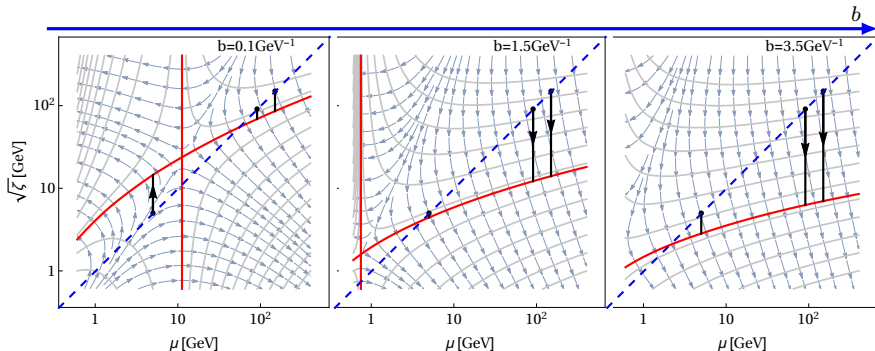
$$\int db e^{iq_T b} C(Q/\mu) F_1(x, b, \mu, \zeta_1) F_1(x, b, \mu, \zeta_2)$$

out: $\mu \sim Q, \zeta_1 = \zeta_2 = Q^2$
in: $\mu, \zeta_1, \zeta_2 \sim ???$



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ζ -prescription



Expression for TMD cross-section is exceptionally simple

$$\frac{d\sigma}{dydQ^2d^2\mathbf{q}_T} = \sigma_0 \int d^2b e^{i(\mathbf{b}\cdot\mathbf{q}_T)} H_{ff'} \left(\frac{Q}{\mu} \right) \left(\frac{Q^2}{\zeta Q} \right)^{-2\mathcal{D}(Q,b)} F_{f\leftarrow h}(x_1, b) F_{f'\leftarrow h}(x_2, b)$$

Each cross-section requires convolution of 3 NP functions

Rapidity
anomalous dimension
(universal)

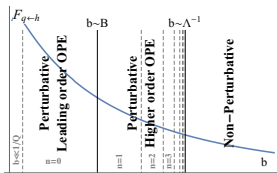
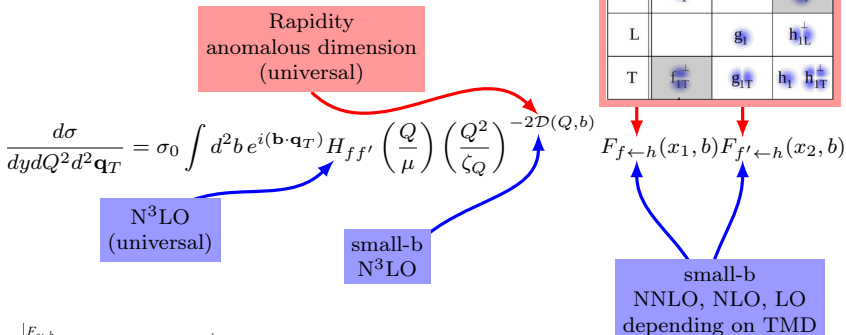
N \ q	U	L	T
U	f_1		h_{1T}
L		g_1	h_{1L}
T	f_{1T}	g_{1T}	h_{1T} h_{TT}

$F_{f \leftarrow h}(x_1, b) F_{f' \leftarrow h}(x_2, b)$

$$\frac{d\sigma}{dy dQ^2 d^2\mathbf{q}_T} = \sigma_0 \int d^2b e^{i(\mathbf{b} \cdot \mathbf{q}_T)} H_{ff'} \left(\frac{Q}{\mu} \right) \left(\frac{Q^2}{\zeta_Q} \right)^{-2\mathcal{D}(Q,b)}$$



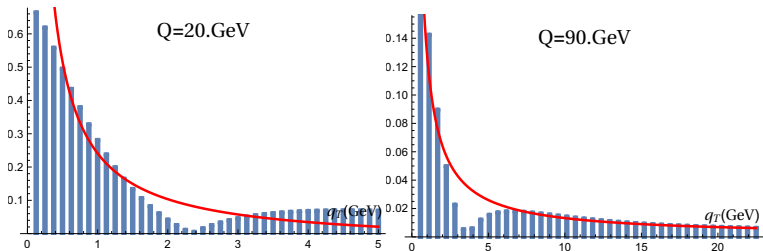
Each cross-section requires convolution of 3 NP functions



Small-b \leftrightarrow large q_T
At small-b everything is perturbative turns to resummation

How large is non-perturbative component?

$DY(\text{resummation})/DY(\text{TMD})$

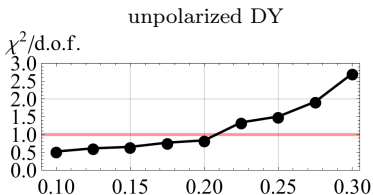


At $q_T \lesssim 0.1Q$ non-perturbative corrections are essential

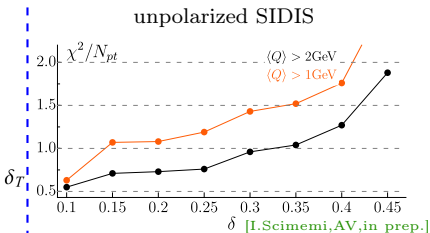


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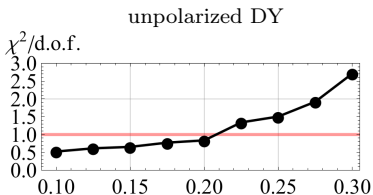
Test by inclusion of the data with
 $q_T < \delta \cdot Q$



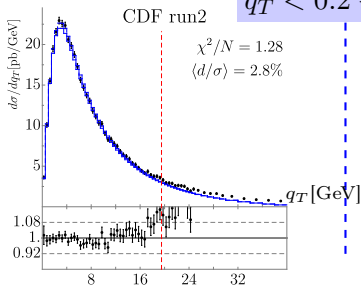
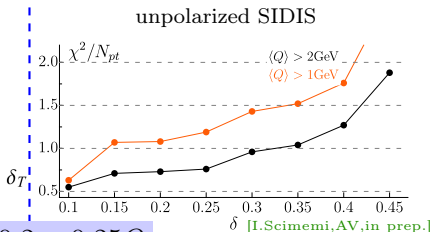
[I.Scimemi,AV,1706.01473]



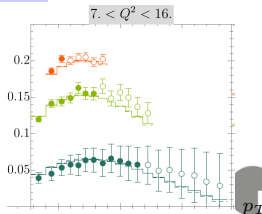
Test by inclusion of the data with
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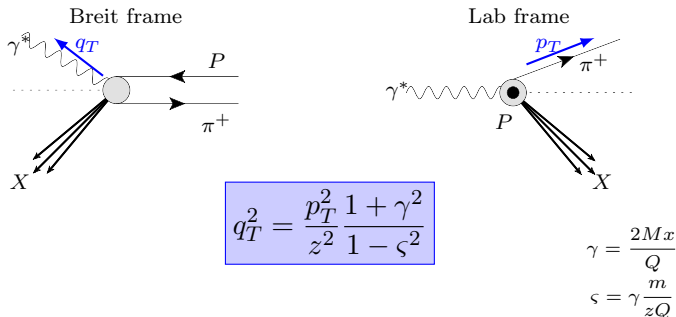
[I.Scimemi, AV, 1706.01473]



$q_T < 0.2 - 0.25Q$



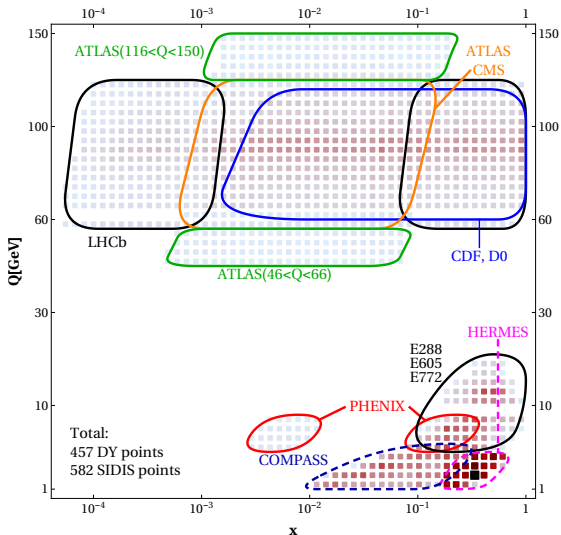
Proof of factorization is done in the Breit frame



In practice: $q_T < 0.25Q$

- ▶ Most part of data is not TMD factorisable.
- ▶ Low z 's are not accessible
- ▶ H1, ZEUS data have no TMD points, too low z .

Global fit of unpolarized DY+SIDIS [I.Scimemi,AV,in prep.]



High energy DY: 194 points
Low energy DY: 263 points
(Low energy) SIDIS: 582 points

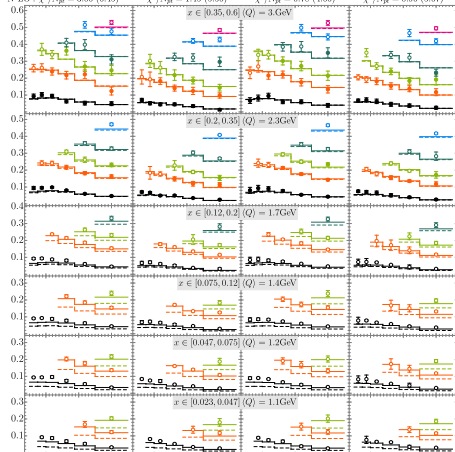
$$\text{NNLO: } \chi_{global}^2 / N_{pt} = 0.95$$

$$\text{N}^3\text{LO: } \chi_{global}^2 / N_{pt} = 1.05$$

artemide v.2.02

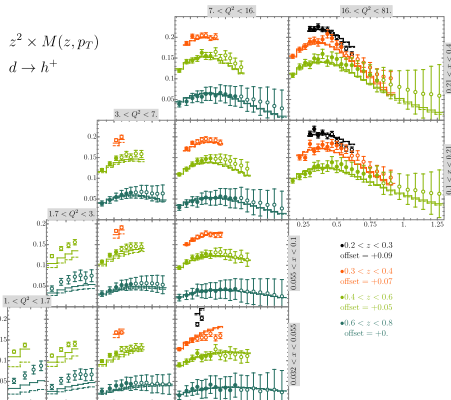


$p \rightarrow \pi^+$ $p \rightarrow \pi^-$ $d \rightarrow \pi^+$ $d \rightarrow \pi^-$
 NNLO: $\chi^2/N_{pt} = 2.20$ (2.64) $\chi^2/N_{pt} = 1.12$ (2.31) $\chi^2/N_{pt} = 0.57$ (1.46) $\chi^2/N_{pt} = 0.74$ (1.91)
 N³LO: $\chi^2/N_{pt} = 3.06$ (6.45) $\chi^2/N_{pt} = 1.45$ (5.56) $\chi^2/N_{pt} = 0.78$ (4.66) $\chi^2/N_{pt} = 0.96$ (5.67)



- $p_T \in [0., 0.15]$ GeV
offset = +0.
- $p_T \in [0.15, 0.25]$ GeV
offset = +0.05
- $p_T \in [0.25, 0.35]$ GeV
offset = +0.1
- $p_T \in [0.35, 0.45]$ GeV
offset = +0.2
- $p_T \in [0.45, 0.6]$ GeV
offset = +0.35
- $p_T \in [0.6, 0.8]$ GeV
offset = +0.45

Example of SIDIS data



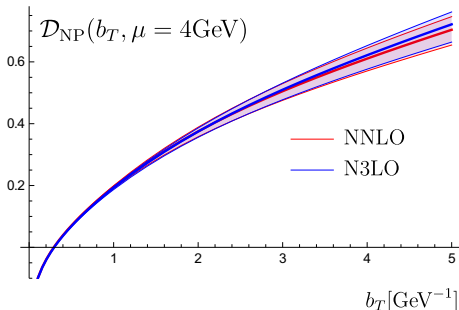
TMD evolution kernel measured in one process suites for all

Using DY extract:
TMDPDF
TMD evolution



Using SIDIS extract:
TMDFF

$$\chi^2/N_{pt}=1.00$$



Using DY+SIDIS extract:
TMDPDF
TMDFF
TMD evolution

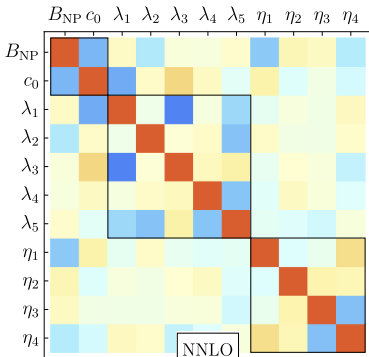
$$\chi^2/N_{pt}=0.95$$

Using π DY extract:
 π TMDPDF

$$\chi^2/N_{pt}=1.3$$



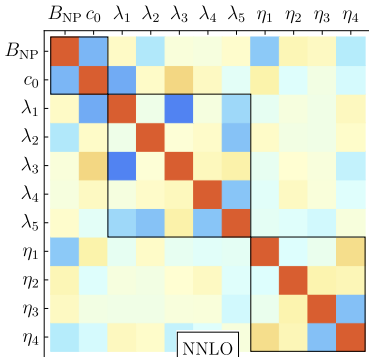
Evolution : 2 parameters
 TMDPDF : 5 parameters + PDF
 TMDFE : 4 parameters + NNFF



Different NP functions are almost decorrelated



Evolution : 2 parameters
 TMDPDF : 5 parameters + PDF
 TMDFF : 4 parameters + NNFF



Different NP functions are almost decorrelated

Fit quality essentially depends on the collinear input.

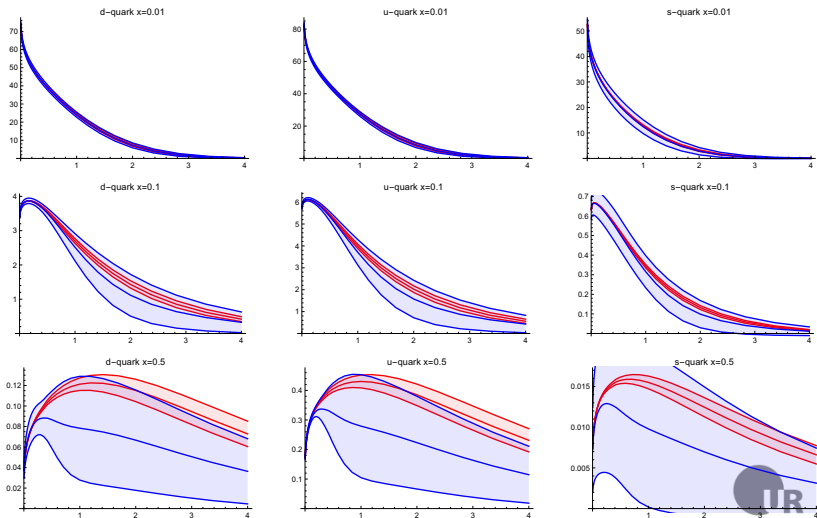
Vary NNPDF within the 1σ band

$$\chi^2/N_{pt} \in [0.8, 35.]$$

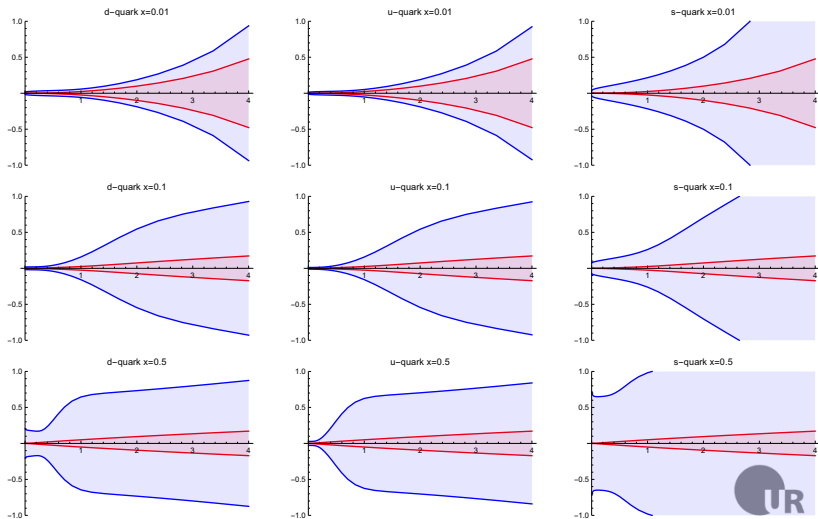
We cannot estimate accurately we the PDF uncertainty.



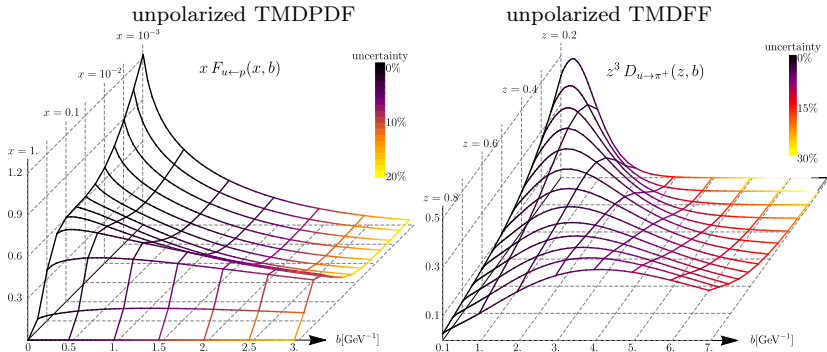
The first attempt to estimate a “true” uncertainty band for TMDPDF



The first attempt to estimate a “true” uncertainty band for TMDPDF



TMD-distributions do not follow perturbative models



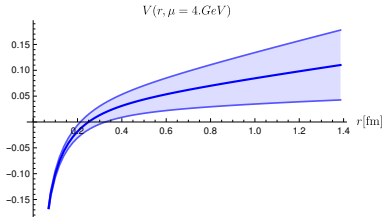
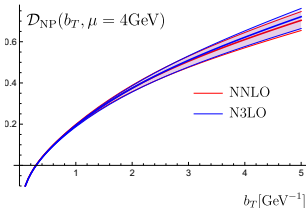
... systematic error 10-20%
... evolution does not have this systematics



TMDs are the window to a QCD dynamics

Example

In stochastic vacuum model



$$V(r) = r \frac{\pi}{4} \mathcal{D}''(0) + \frac{\mathcal{D}'(0)}{2} + \frac{r^2}{2} \int_r^\infty \frac{dx}{x^2} \frac{\mathcal{D}'(x)}{\sqrt{x^2 - r^2}}$$

QCD string tension: $\sigma = 0.014 \pm 0.01 \text{ GeV}^2$ vs. $\sim 0.1 - 0.2 \text{ GeV}^2$



Conclusion

TMD distributions are known, and predictions can be generated with it

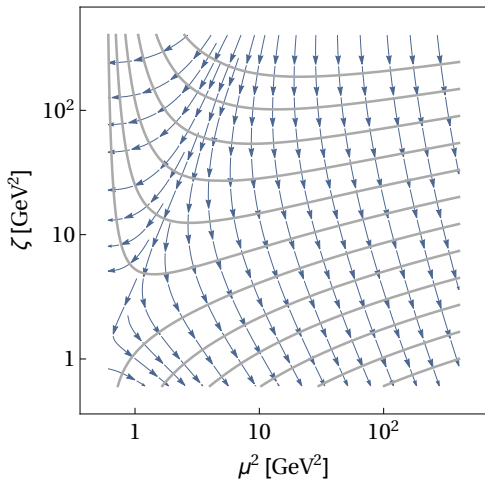
But there are plenty of unknowns...

- ▶ Flavor separation...
- ▶ Collinear parts ...
- ▶ Power corrections ...
- ▶ ...

artemide

- ▶ Fortran 95
- ▶ Python interface (**harpy**)
- ▶ Modular structure (from NP-input to cross-section) + LHAPDF interface
- ▶ Variety of theory tools and implementations
- ▶ Growing functionality

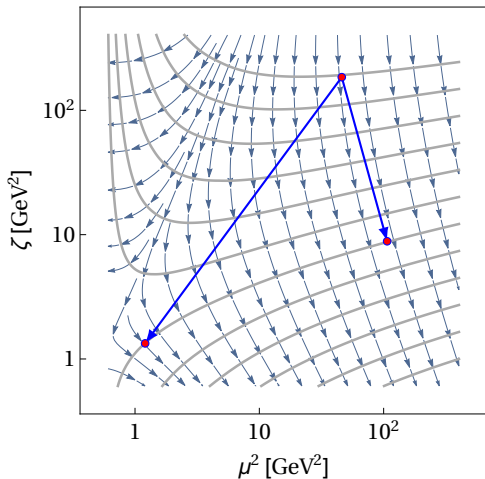
TMD distribution is not defined by a scale (μ, ζ)
It is defined by an equipotential line.



The scaling is defined by
~~a difference between scales~~
a difference between potentials



TMD distribution is not defined by a scale (μ, ζ)
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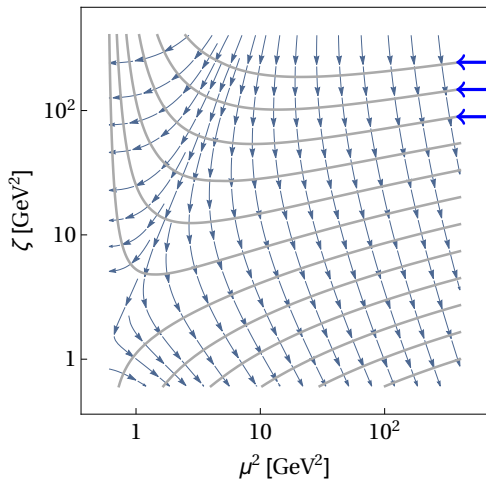


The scaling is defined by
~~a difference between scales~~
a difference between potentials

Evolution factor to both points
is the same
although the scales are
different by 10^2GeV^2



TMD distributions on the same equipotential line are equivalent.



We can enumerate them by a lines
not by (μ, ζ)

$$F(x, b; \mu, \zeta) \rightarrow F(z, b; \text{line})$$

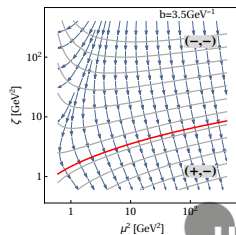
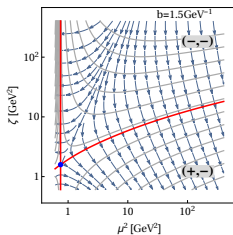
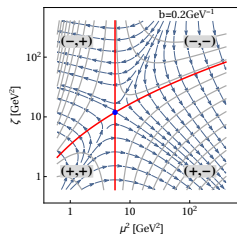


There is a unique line which passes through all μ 's

The optimal TMD distribution

$$F(x, b) = F(x, b; \mu, \zeta_\mu)$$

where ζ_μ is the special line.



Comparison with other evolutions

