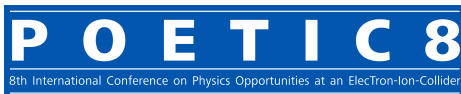


# Transverse momentum dependent (TMD) factorization in perturbation theory

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TMD factorization is a complicated composition of perturbative and non-perturbative functions.

$$\frac{d\sigma}{dX} \simeq \int db e^{ibq_T} H(Q) \{R(Q \rightarrow (\mu_i, \zeta_i))\}^2 F_1(x, b, \mu_i, \zeta_i) F_2(x, b, \mu_i, \zeta_i)$$



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Hard coef.  
Perturbative



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UV anomalous dimension  
Perturbative

Evolution factor  
 $\gamma_F$   $\mathcal{D}(b)$

Rap. anomalous dimension  
 $b \ll \Lambda^{-1}$  Perturbative  
 $b \gtrsim \Lambda^{-1}$  Non-Perturbative



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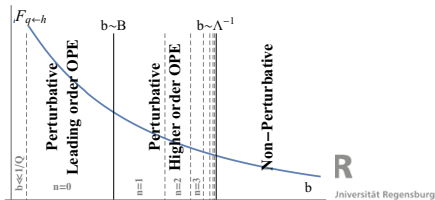
UV anomalous dimension  
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Evolution factor  
 $\gamma_F$   $D(b)$

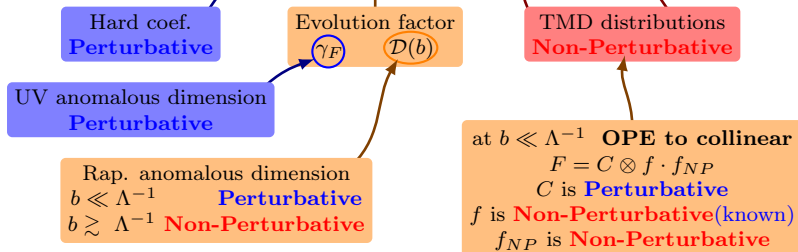
TMD distributions  
Non-Perturbative

at  $b \ll \Lambda^{-1}$  OPE to collinear  
 $F = C \otimes f \cdot f_{NP}$   
 $C$  is Perturbative  
 $f$  is Non-Perturbative (known)  
 $f_{NP}$  is Non-Perturbative



TMD factorization is a complicated composition of perturbative and non-perturbative functions.

$$\frac{d\sigma}{dX} \simeq \int db e^{ibq_T} H(Q) \{R(Q \rightarrow (\mu_i, \zeta_i))\}^2 F_1(x, b, \mu_i, \zeta_i) F_2(x, b, \mu_i, \zeta_i)$$

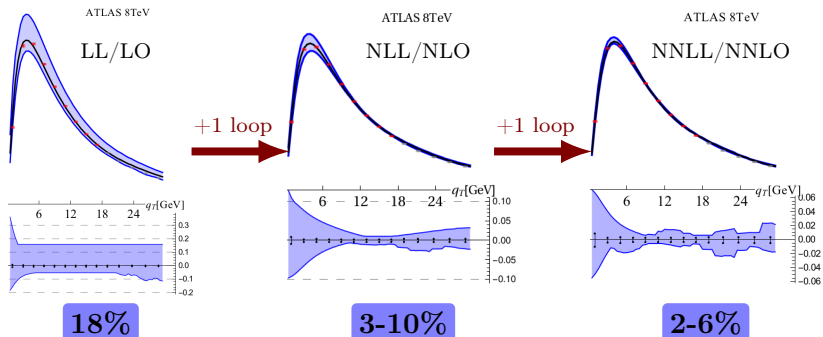


Balance

	<b>Perturbative</b>	<b>Non-Perturbative</b>	
H	Hard coeff.	collinear distributions	$f, d, ..$
$\gamma$	UV anomalous dimension	TMD behaviour	$f_{NP}$
$\mathcal{D}$	rap. anomalous dimension	rap. anomalous dimension	$\mathcal{D}_{NP}$
$C$	mathing coefficients		

Perturbative input is essential for precise phenomenology.  
All perturbative inputs are important.

## High energy

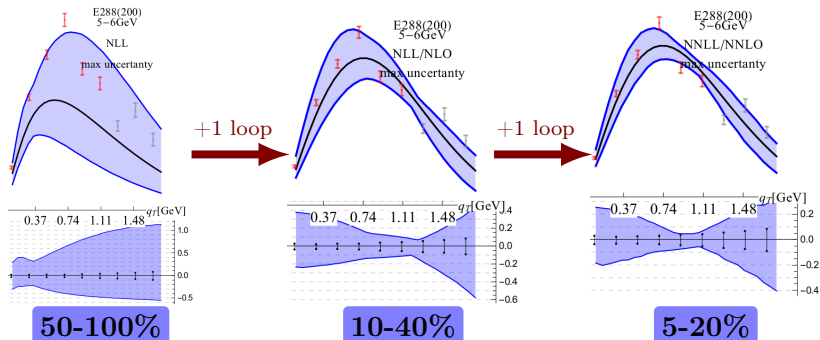


plots from [I.Scimemi,AV; 1706.01473]



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## Low energy



plots from [I.Scimemi,AV; 1706.01473]

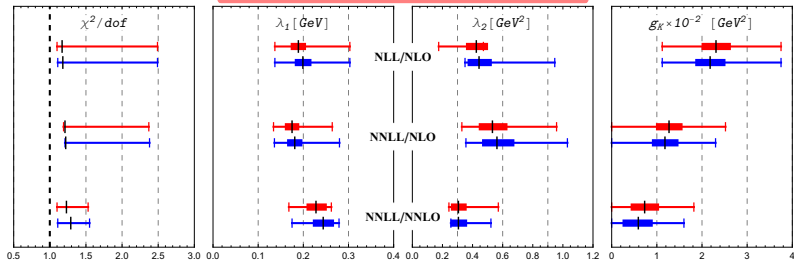




Perturbative input is essential for precise phenomenology.  
All perturbative inputs are important.

## Effect on the TMD extraction

LL/LO is not shown (errors too large)



plots from [I.Scimemi, AV; 1706.01473]

## Conclusion

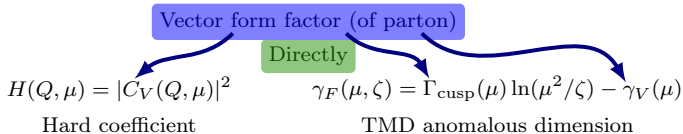
Perturbation input is important at both low and high energies.  
The higher order – the better.

There was significant progress in the perturbative calculus

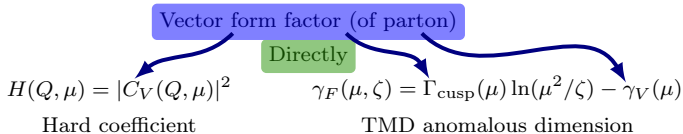
- Schemes of perturbative calculation were consistently formulated
- 3-loop TMD evolution [Li,Zhu,1604.01404][AV,1610.05791]
- All TMD distributions (tw-2) were recalculated in the same scheme [Scimemi,AV,1702.06558][Buffing,Diehl,Kasemets,1708.03528].
- (Rapidity divergences part) of TMD factorization was proven [AV,1707.07606].
- More to come soon.

	known	in this talk
H (universal)	3-loop [2010]	
$\gamma_F$ (universal)	3-loop [2010]	
$\mathcal{D}$ (universal)	3-loop [2016]	✓
$f_1$ (unpol. PDF)	2-loop [2015]	
$d_1$ (unpol.FF)	2-loop [2016]	
$g_1$ (helicity)	1-loop [2013]	
$h_1$ (transvercity)	1-loop [2013]	✓
$h_{1T}^\perp$ (pretzelocity)	1-loop (=0) [2017]	✓
$f_{1T}^\perp$ (Sivers)	?tree [2014]	✓
$d_{1T}^\perp$ (Collins)	unknown	✓
$h_1^\perp$ (Boer-Mulders)	unknown	✓
$g_{1T}, h_{1L}$ (worm-gear T,L)	unknown	✓

Not everything requires (full) calculation.  
 Universal quantities could be taken from outside.



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Universal quantities could be taken from outside.



Soft/rapidity anomalous dimension correspondance

Rapidity anomalous dimension

$$\mathcal{D}(\mu, b) \quad (= -\bar{K}/2)$$

TMD physics



Soft anomalous dimension

$$\gamma_{\text{soft}}(\mu, v_{ij})$$

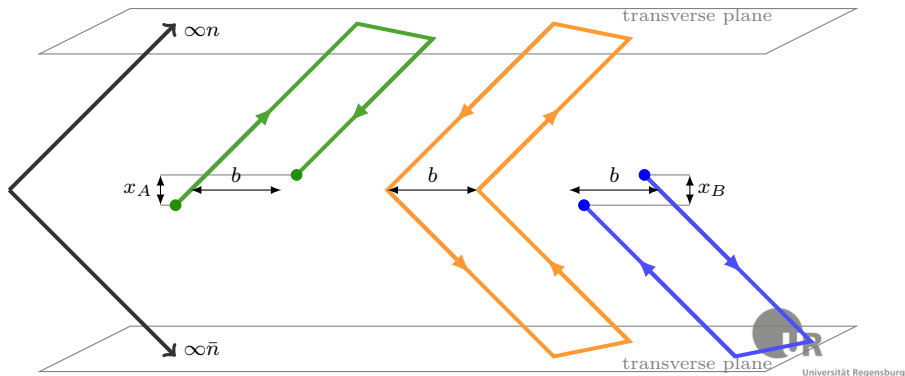
jet production

Byproduct of rapidity divergences renormalization theorem  
that completes the TMD factorization



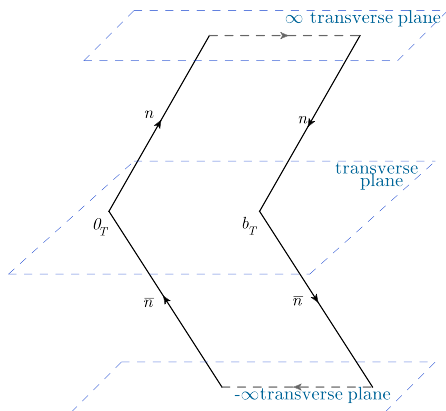
TMD factorization is two successive factorizations  
**The first:** collinear factorization [Collins' book], [SCET]  
**The second:** rapidity divergences factorization [AV,1707.07606]

$$\frac{d\sigma}{dX} \simeq H(Q) \int \frac{d^2b}{(2\pi)^2} e^{i(bk)_T} f(x_A, b) S(b) f(x_B, b)$$



Soft factor is vacuum matrix element of a Wilson loop.

$$S(b) = \langle 0 | [0, \pm\infty n] [\pm\infty n + \mathbf{b}, \mathbf{b}] [\mathbf{b}, \pm\infty \bar{n}] [\pm\infty \bar{n}, 0] | 0 \rangle$$

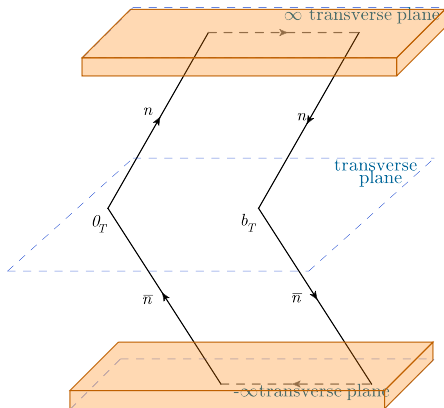


Various kinds of divergences



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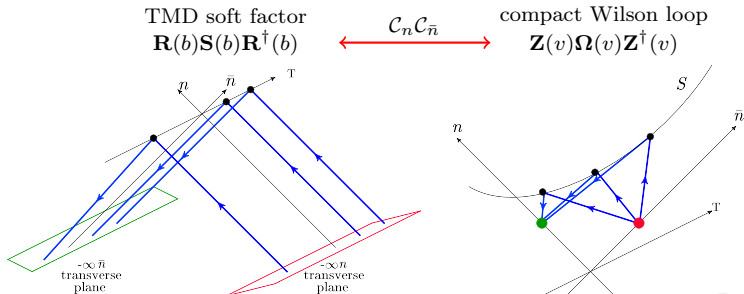
### Various kinds of divergences

- Ultraviolet (renormalized)
- Collinear (cancel in sum)
- Mass divergences (cancel in sum)
- Rapidity divergences (???)

There is a geometrical transformation  
that turns rapidity divergence to UV

$$\mathcal{C}_{\bar{n}} : \{x^+, x^-, x_\perp\} \rightarrow \left\{ \frac{-1}{2a} \frac{1}{\lambda + 2ax^+}, x^- + \frac{ax_\perp^2}{\lambda + 2ax^+}, \frac{x_\perp}{\lambda + 2ax^+} \right\}$$

- The UV renormalization imposes rapidity divergence renormalization

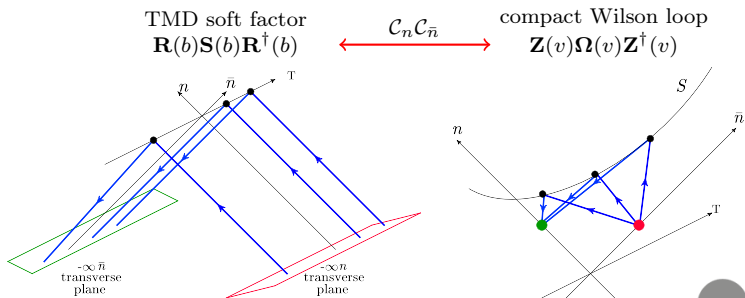




## Rapidity divergence renormalization theorem

Rapidity divergences for TMD soft factor can be renormalized for each direction separately (also holds for MPS).

- Immediate in conformal field theory
- Proof by iteration in QCD (1-loop in conformal)



## Soft/Rapidity anomalous dimension correspondence

Exact relation

$$\gamma_{\text{soft}}(v) = 2\mathcal{D}(\mathbf{b}, \epsilon^*)$$

How to use it?

- Physical value is  $\mathbf{D}(\{\mathbf{b}\}, 0)$
- $\epsilon^* = 0 - a_s\beta_0 - a_s^2\beta_1 - a_s^3\beta_2 - \dots$
- We can compare order by order in PT

$$\mathbf{D}_1(\{b\}) = \frac{1}{2}\gamma_1(\{v\}),$$

$$\mathbf{D}_2(\{b\}) = \frac{1}{2}\gamma_2(\{v\}) + \beta_0\mathbf{D}'_1(\{b\}),$$

$$\mathbf{D}_3(\{b\}) = \frac{1}{2}\gamma_3(\{v\}) + \beta_0\mathbf{D}'_2(\{b\}) + \beta_1\mathbf{D}'_1(\{b\}) - \frac{\beta_0^2}{2}\mathbf{D}''_1(\{b\}),$$

Always gain one order!

This simple exercise gives 3-loop rapidity anomalous dimension.

$$\begin{aligned}
 \mathcal{D}_{L=0}^{(3)} = & -\frac{C_A^2}{2} \left( \frac{12328}{27} \zeta_3 - \frac{88}{3} \zeta_2 \zeta_3 - 192 \zeta_5 - \frac{297029}{729} + \frac{6392}{81} \zeta_2 + \frac{154}{3} \zeta_4 \right) \\
 & - \frac{C_A N_f}{2} \left( -\frac{904}{27} \zeta_3 + \frac{62626}{729} - \frac{824}{81} \zeta_2 + \frac{20}{3} \zeta_4 \right) - \\
 & \frac{C_F N_f}{2} \left( -\frac{304}{9} \zeta_3 + \frac{1711}{27} - 16 \zeta_4 \right) - \frac{N_f^2}{2} \left( -\frac{32}{9} \zeta_3 - \frac{1856}{729} \right)
 \end{aligned}$$

- Coincides with the one calculated directly [\[Li,Zhu,1604.01404\]](#)

Unfortunately, there is something to calculate

N \ q	U	L	T
U	$f_1$		$h_1^\perp$
L		$g_1$	$h_{1L}^\perp$
T	$f_{1T}^\perp$	$g_{1T}$	$h_1$ $h_{1T}^\perp$

OPE at small values of  $b$

$$F(x, b) = C(x, b; \mu_{\text{OPE}}) \otimes f(x, \mu_{\text{OPE}}) + \dots$$

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T	$f_{1T}^\perp$	$g_{1T}$	$h_1$ $h_{1T}^\perp$

○ Leading twist-2

- $f_1 \rightarrow$  Unpol.PDF (2-loop)  
[1403.6451,1604.07869]
- $d_1 \rightarrow$  Unpol.FF (2-loop)  
[1509.06392,1604.07869]
- $g_1 \rightarrow$  helicity PDF (1-loop)  
[1303.2129,1702.06558,1708.03528]
- $h_1 \rightarrow$  transverscity PDF (1-loop)  
[1303.2129,1702.06558,1708.03528]
- $h_{1T}^\perp \rightarrow$  transverscity PDF(1-loop)  
(=0) [1702.06558]

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T	$f_{1T}^\perp$	$g_{1T}$	$h_1$ $h_{1T}^\perp$

○ Leading twist-2

○ Leading twist-3

- $f_{1T}^\perp \rightarrow$  Qui-Sterman  
all calculations made for cross-section
- $h_1^\perp \rightarrow ??$

These function must change sign under  $DY \leftrightarrow \text{SIDIS}$

OPE at small values of  $b$

$$F(x, b) = C(x, b; \mu_{\text{OPE}}) \otimes f(x, \mu_{\text{OPE}}) + \dots$$



Unfortunately, there is something to calculate

N \ q	U	L	T
U	$f_1$		$h_1$
L		$g_1$	$h_{1L}$
T	$f_{1T}$	$g_{1T}$	$h_1$ $h_{1T}$

○ Leading twist-2

○ Leading twist-3

○ Leading twist-2 and 3

- $h_{1L}^\perp \rightarrow ?? + WW$  (helicity PDF)
- $g_{1T} \rightarrow ?? + WW$  (transverscity PDF)

WW part can be found in  
[Kanazawa, et al; 1512.07233]

OPE at small values of  $b$

$$F(x, b) = C(x, b; \mu_{\text{OPE}}) \otimes f(x, \mu_{\text{OPE}}) + \dots$$

## New 2-loop results

$$\Phi_{q \leftarrow h}^{[i\sigma^{\alpha+}\gamma_5]}(x, \mathbf{b}) = s_T^\alpha h_1(x, \mathbf{b}) + 2 \left( \frac{g_T^{\alpha\mu}}{2} + \frac{b^\alpha b^\mu}{\mathbf{b}^2} \right) s_{T\mu} h_{1T}^\perp(x, \mathbf{b}) + b_\mu(\dots)^{\alpha\mu}$$





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transvercity TMDPDF



transvercity PDF

$$h_1(x, \mathbf{b}) = C \otimes \delta f(x) + \mathbf{b}^2 \dots$$

tree+1loop



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tree+1loop

linear in  $\mathbf{b}$



twist-3



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transvercity TMDPDF



transvercity PDF

$$h_1(x, \mathbf{b}) = C \otimes \delta f(x) + \mathbf{b}^2 \dots$$

tree+1loop

pretzelocity TMDPDF



transvercity PDF

$$h_{1T}^\perp(x, \mathbf{b}) = C_T^\perp \otimes \delta f(x) + \mathbf{b}^2 \dots$$

linear in  $\mathbf{b}$ 

twist-3

$$C_T^\perp(x, \mathbf{b}) = -4a_s^2 \left[ 4 \left( C_F^2 - \frac{C_F C_A}{2} \right) \left( \bar{x} \ln \bar{x} + x \ln x - \frac{3}{2} \bar{x} \right) - C_F^2 \bar{x} (x - \dots) \right]$$

Preliminary



## New 2-loop results

$$\Phi_{q \leftarrow h}^{[i\sigma^{\alpha+}\gamma_5]}(x, \mathbf{b}) = s_T^{\alpha} h_1(x, \mathbf{b}) + 2 \left( \frac{g_T^{\alpha\mu}}{2} + \frac{b^\alpha b^\mu}{\mathbf{b}^2} \right) s_{T\mu} h_{1T}^\perp(x, \mathbf{b}) + b_\mu(\dots)^{\alpha\mu}$$

transvercity TMDPDF



transvercity PDF

$$h_1(x, \mathbf{b}) = C \otimes \delta f(x) + \mathbf{b}^2 \dots$$

tree+1loop+2loop

pretzelocity TMDPDF



transvercity PDF

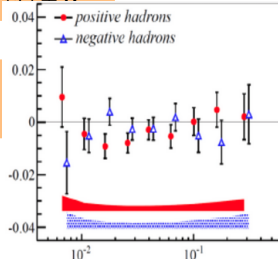
$$h_{1T}^\perp(x, \mathbf{b}) = C_T^\perp \otimes \delta f(x) + \mathbf{b}^2$$

linear in  $\mathbf{b}$ 

twist-3

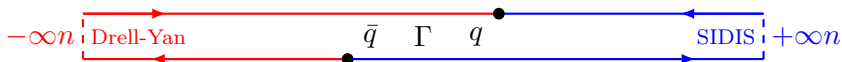
$$C_T^\perp(x, \mathbf{b}) = -4a_s^2 \left[ 4 \left( C_F^2 - \frac{C_F C_A}{2} \right) \left( \bar{x} \ln \bar{x} + x \ln x - \right. \right.$$

The first model independent result for  
pretzelocity  
Natural explanation of smallness of  
 $\sin(3\varphi - \varphi_s)$  asymmetry



Leading matching of all polarized distributions at twist-3

$$U_{\text{DY}}^{\Gamma}(z, \mathbf{b}) = \bar{q}(zn + \mathbf{b})[zn + \mathbf{b}, -\infty n + \mathbf{b}]\Gamma[-\infty n - \mathbf{b}, -zn - \mathbf{b}]q(-zn - \mathbf{b}),$$

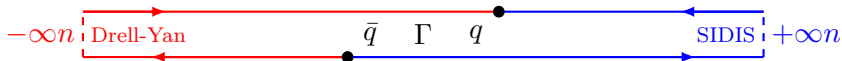


$$U_{\text{SIDIS}}^{\Gamma}(z, \mathbf{b}) = \bar{q}(zn + \mathbf{b})[zn + \mathbf{b}, +\infty n + \mathbf{b}]\Gamma[+\infty n - \mathbf{b}, -zn - \mathbf{b}]q(-zn - \mathbf{b}),$$



## Leading matching of all polarized distributions at twist-3

$$U_{\text{DY}}^{\Gamma}(z, \mathbf{b}) = \bar{q}(zn + \mathbf{b})[zn + \mathbf{b}, -\infty n + \mathbf{b}]\Gamma[-\infty n - \mathbf{b}, -zn - \mathbf{b}]q(-zn - \mathbf{b}),$$



$$U_{\text{SIDIS}}^{\Gamma}(z, \mathbf{b}) = \bar{q}(zn + \mathbf{b})[zn + \mathbf{b}, +\infty n + \mathbf{b}]\Gamma[+\infty n - \mathbf{b}, -zn - \mathbf{b}]q(-zn - \mathbf{b}),$$

$$U_{\text{DY(SIDIS)}}^{\Gamma}(z, \mathbf{b}) = O^{\Gamma}(z) + \mathbf{b}^{\mu} \left\{ \lim_{x \rightarrow z} \frac{\partial}{\partial x^{\mu}} O^{\Gamma}(x) - i \int_{-1}^1 dv v z T_{\mu}^{\Gamma}(z, v z, -z) \right. \\ \left. + (-) i \left( \int_{-\infty}^z + \int_{-\infty}^{-z} \right) d\tau T_{\mu}^{\Gamma}(z, \tau, -z) \right\} + \mathcal{O}(\mathbf{b}^2),$$

$$O^{\Gamma}(z) = \bar{q}(z)\Gamma[z, -z]q(z),$$

$$T_{\mu}^{\Gamma}(z, \tau, -z) = \bar{q}(z)[z, \tau]\Gamma F_{\mu+}(\tau)[\tau, -z]q(-z)$$

## Leading matching of all polarized distributions at twist-3

Collins  $f_{1T}^\perp(x, \mathbf{b}) = +(-)\pi T(x, x) + O(\mathbf{b}^2)$

Boer-Mulders  $h_1^\perp(x, \mathbf{b}) = -(+)\pi \delta T_\epsilon(x, x) + O(\mathbf{b}^2)$

Worm-gear T  $g_{1T}^\perp(x, \mathbf{b}) = 2 \int \frac{dx_2}{x_2^2} \int_{x-x_2}^x du \Delta T(u, u+x_2)$   
 $-\frac{1}{2} \int du \frac{u}{|u|} (g_1(u) - g_T(u)) + O(\mathbf{b}^2)$

Worm-gear L  $h_{1L}^\perp(x, \mathbf{b}) = -2 \int \frac{dx_2}{x_2^2} \int_{x-x_2}^x du \delta T_g(u, u+x_2)$   
 $+\frac{1}{2} \int du \frac{u}{|u|} (h_1(u) - h_L(u)) + O(\mathbf{b}^2)$

### Qiu-Sterman functions

$$\langle \mathcal{T}_\mu^{\gamma^+} \rangle \sim \tilde{s}_{\mu T} T \quad \langle \mathcal{T}_\mu^{\gamma^+ \gamma^5} \rangle \sim s_{\mu T} \Delta T \quad \langle \mathcal{T}_\mu^{\sigma^{\alpha+} \gamma^5} \rangle \sim \epsilon_T^{\mu\alpha} \delta T_\epsilon + g_T^{\mu\alpha} \delta T_g$$

Functions  $g_T$  and  $h_L$  are compositions of twist-2/3 PDFs.

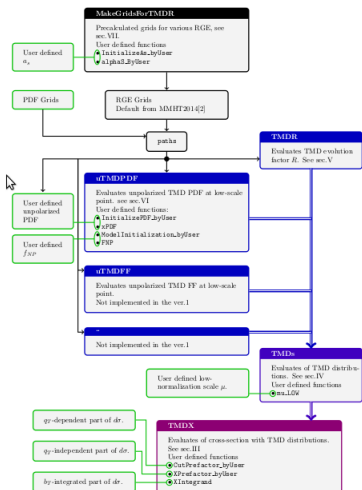
# Conclusion

- Factorization of rapidity divergences is proven (for SIDIS assuming universality)
- TMD evolution is complete 3-loop
- Small-b matching is known for all TMDs (see table)

	known	soon to come
H (universal)	3-loop [2010]	
$\gamma_F$ (universal)	3-loop [2010]	
$\mathcal{D}$ (universal)	3-loop [2016]	
$f_1$ (unpol. PDF)	2-loop [2015]	
$d_1$ (unpol.FF)	2-loop [2016]	
$g_1$ (helicity)	1-loop [2013]	
$h_1$ (transverscity)	1-loop [2013]	<b>2-loop</b>
$h_{1T}^\perp$ (pretzelocity)	1-loop (=0) [2017]	<b>2-loop</b>
$f_{1T}^\perp$ (Sivers)	?tree [2014]	<b>tree</b>
$d_{1T}^\perp$ (Collins)	unknown	<b>tree</b>
$h_{1T}^\perp$ (Boer-Mulders)	unknown	<b>tree</b>
$g_{1T}, h_{1L}$ (worm-gear T,L)	unknown	<b>tree</b>







Tool for phenomenological studies of TMDs with high-end theory input

- FORTRAN 90 code
- Module structure
- Convolutions, evolution (LO, NLO, NNLO)
- Fourier to  $q_T$ -space, integrations over phase space
- Scale-variation ( $\zeta$ -prescription)
- User defined PDFs, scales,  $f_{NP}$
- Efficient code ( $\sim 10^9$  TMDs  $\sim 6$ . min at NNLO)

Currently ver 1.1 (soon 1.2 with TMD FFs)

Available at: <https://teorica.fis.ucm.es/artemide>

